## Self-avoiding walks in wedges

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# Self-avoiding walks in wedges 

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#### Abstract

We consider the number of self-avoiding walks confined to a subset $\mathbb{Z}^{d}(f)$ of the $d$-dimensional hypercubic lattice $\mathbb{Z}^{d}$, such that the coordinates $\left(x_{1}, x_{2}, \ldots, x_{d}\right)$ of each vertex in the walk satisfy $x_{1} \geqslant 0$ and $0 \leqslant x_{k} \leqslant f_{k}\left(x_{1}\right)$ for $k=2,3, \ldots, d$. We show that if $f_{k}(x) \rightarrow \infty$ as $x \rightarrow \infty$, the connective constant of walks in $\mathbb{Z}^{d}(f)$ is identical to the connective constant of walks in $\mathbb{Z}^{d}$. We also explore conditions on $f_{k}$ which lead to a smaller connective constant for walks in $\mathbb{Z}^{d}(f)$ and, in particular, consider walks between two parallel ( $d-1$ )-dimensional hyperplanes. Finally we contrast some of these results with recent work by Grimmett on percolation on subsets of the square lattice.


## 1. Introduction

Self-avoiding walks with various geometrical constraints have been studied as models of polymer adsorption (Hammersley et al 1982), steric stabilisation of dispersions (Dolan and Edwards 1975, Levine et al 1978) and the behaviour of polymers in slits and tubes (Daoud and de Gennes 1977, Wall et al 1977, Wall and Klein 1979). Because of the analogy between the polymer problem and the behaviour of magnetic systems (Daoud et al 1975) those models are also useful in understanding surface magnetism and the behaviour of slabs of magnetic material (Binder 1983).

Cardy (1983) has recently considered the critical behaviour of a magnet in $d$ dimensions, bounded by two ( $d-1$ )-dimensional hyperplanes which meet at an edge. In particular he used renormalisation group methods to study the critical exponents for correlation functions and susceptibilities for spins close to the edge. The polymer analogue of this problem can be obtained by considering $D$-dimensional spins on each $d$-dimensional lattice site and letting $D \rightarrow 0$, (Daoud et al 1975, Barber et al 1978). One is then interested in the numbers of self-avoiding walks starting at the origin and confined to lie between these planes, and the subsets of these walks which either return to one of these planes, or to the edge, at their last step.

In this paper we shall prove that these three classes of self-avoiding walks have the same connective constant ( $\kappa$ ) as self-avoiding walks with no such geometrical restrictions. In fact we study the more general problem of walks on a $d$-dimensional hypercubic lattice such that the coordinates ( $x_{1}, x_{2}, \ldots, x_{d}$ ) of each vertex in the walk satisfy $0 \leqslant x_{1}, 0 \leqslant x_{k} \leqslant f_{k}\left(x_{1}\right)$ for $2 \leqslant k \leqslant d$. We establish sufficient conditions on the $f_{k}$ to ensure that the walks, subject to these restrictions, have connective constant $\kappa$. We also discuss conditions on $f_{k}$ which are sufficient to ensure that the connective constant is strictly less than $\kappa$ and, in particular, study walks confined between two parallel

[^0]planes, a distance $L$ apart. In this case we show that the connective constant $\kappa(L)$ is a strictly monotone function of $L$ and tends to $\kappa$ as $L$ tends to infinity.

Finally, we compare some of our results with recent work by Grimmett (1981, 1983) on percolation on subsets of the square lattice. The behaviour for percolation is characteristically different to that for self-avoiding walks.

## 2. Definitions and notation

We work in $d$-dimensional Euclidean space $\mathbb{R}^{d}$ with $d \geqslant 2$, writing $x=\left(x_{1}, x_{2}, \ldots, x_{d}\right)$ for a typical point and its coordinates: here bold-face type signifies vectors, so that a typical sequence of points $x_{1}, x_{2}, \ldots$ should not be confused with the coordinates $x_{1}, x_{2}, \ldots$ of $x$. The vector inequality

$$
\begin{equation*}
a \leqslant x \leqslant b \tag{2.1}
\end{equation*}
$$

is an abbreviation for

$$
\begin{equation*}
a_{k} \leqslant x_{k} \leqslant b_{k} \quad(k=1,2, \ldots, d) \tag{2.2}
\end{equation*}
$$

and the set of all $\boldsymbol{x}$ lying in a closed interval of the form (2.1) is called a box. A box has $2 d$ faces, these being the $(d-1)$-dimensional subsets obtained by substituting equality for one of the $2 d$ inequalities in (2.2); but not all of these faces will be distinct unless $\boldsymbol{a}<\boldsymbol{b}$. The smallest box that contains some bounded set of points $S$ is called the snug box for $S$. A translate of $S$ is the image of $S$ under a transformation $\boldsymbol{x} \rightarrow \boldsymbol{x}+\boldsymbol{a}$ for some constant vector $a$. We write $e_{1}, \ldots, e_{d}$ for the unit vectors in the positive directions of the coordinate axes $\mathbb{R}^{d}$.

Given a set of non-negative functions $f_{2}, f_{3}, \ldots, f_{d}$, the set of all points $x$ whose coordinates satisfy

$$
\begin{equation*}
x_{1} \geqslant 0, \quad f_{k}\left(x_{1}\right) \geqslant x_{k} \geqslant 0 \quad(2 \leqslant k \leqslant d) \tag{2.3}
\end{equation*}
$$

is called an $f$-wedge and denoted by $\mathbb{R}^{d}(f)$. The hypercubic lattice $\mathbb{Z}^{d}$ is the set of all points of $\mathbb{R}^{d}$ with integer coordinates; and $\mathbb{Z}^{d}(f)$ denotes the intersection of $\mathbb{Z}^{d}$ and $\mathbb{R}^{d}(f)$.

An $n$-stepped walk is an ordered sequence $z_{0}, z_{1}, \ldots, z_{n}$ of points of $\mathbb{Z}^{d}$ such that successive points are unit distance apart; and it is an $n$-stepped self-avoiding walk (abbreviated hereafter as an $n$-SAW) if these $n+1$ points are all distinct. We call the points $z_{0}, z_{1}, \ldots, z_{n}$ the vertices of the walk, we say that a walk $w$ visits a point $z$ of $\mathbb{Z}^{d}$ if $z$ is one of its vertices, and we use adjectives such as first, earliest, last, etc with natural reference to the order of the sequence $z_{0}, \ldots, z_{n}$. Two walks are equivalent if one is a translate of the other; and often we are only interested in the properties of $n$-SAws to within equivalence. In that case we can standardise an $n$-saw $w$ by taking its first vertex at the origin ( $z_{0}=0$ ) and representing it by a sequence of steps $w=$ $\sigma_{1} \sigma_{2} \ldots \sigma_{n}$, where each step $\sigma_{i}=z_{i}-z_{i-1}$ takes one of the $2 d$ possible values $\pm \boldsymbol{e}_{k}$ $(k=1,2, \ldots, d)$. We write $c_{n}$ for the number of $n$-SAWs with $z_{0}=e$ and $C(n)$ for the set of these walks; and we know (Hammersley 1957) that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} n^{-1} \log c_{n}=\kappa, \tag{2.4}
\end{equation*}
$$

where $\kappa$, which depends only on $d$, is the so-called connective constant of $\mathbb{Z}^{d}$. (Note, however, that the term connective constant is sometimes used instead for $\mu=\mathrm{e}^{\kappa}$; so
care is needed to identify the usage in any particular paper in the literature.) Similarly we write $c_{n}(f)$ for the number of $n$-saws with $\boldsymbol{z}_{0}=0$ with all their vertices in $\mathbb{Z}^{d}(f)$. We are then interested in conditions upon $f$ such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} n^{-1} \log c_{n}(f)=\kappa \tag{2.5}
\end{equation*}
$$

may be true or false. We shall see, in due course, that a sufficient condition for the truth of (2.5) is

$$
\begin{equation*}
\lim _{x \rightarrow \infty} f_{k}(x)=\infty \quad(2 \leqslant k \leqslant d) \tag{2.6}
\end{equation*}
$$

and conversely that (2.5) is false if $f_{k}$ is a bounded function of $x$ for at least one value of $k$.

An $n$-stepped self-avoiding circuit (abbreviated as $n$-SAC) is an ( $n-1$ )-SAW whose first and last vertices $z_{0}$ and $z_{n-1}$ are unit distance apart. If $\boldsymbol{z}_{0}, \boldsymbol{z}_{1}, \ldots, \boldsymbol{z}_{n-1}$ are the successive vertices of an $n$-sAC, then any cyclic permutation of these vertices is also an $n$-SAC; and so too is the reverse permutation $z_{n-1}, z_{n-2}, \ldots, z_{0}$ and all the cyclic permutations of this reverse permutation. The resulting set of $2 n n$-SACS that originate from a given $n$-SAC may be regarded as a single geometrical entity called an $n$-stepped self-avoiding polygon (abbreviated as $n$-SAP). (An $n$-SAC is a rooted, directed $n$-SAP.) Two $n$-saps are equivalent if one is a translate of the other. Thus a standardised representative of an equivalence class of $n$-SAPS is an $n$-SAP that visits the origin of $\mathbb{Z}^{d}$, because each $n$-sap is an unrooted undirected cycle. We write $p_{n}$ for the number of inequivalent $n$-saps visiting the origin, and $p_{n}(f)$ for the number of those having all their vertices in $\mathbb{Z}^{d}(f)$. Parity considerations on $\mathbb{Z}^{d}$ require that any $n$-sap must have $n$ even; and hereafter we tacitly adopt the convention that $n$ is always even in any statement involving $n$-SAPs. We know (Hammersley 1961) that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} n^{-1} \log p_{n}=\kappa \tag{2.7}
\end{equation*}
$$

and so we enquire whether or not

$$
\begin{equation*}
\lim _{n \rightarrow \infty} n^{-1} \log p_{n}(f)=\kappa \tag{2.8}
\end{equation*}
$$

If there is some integer $x_{0} \geqslant 0$ such that $f_{k}\left(x_{0}\right)=0$ for all $k=2,3, \ldots, d$, there can be no $n$-SAPS in $\mathbb{Z}^{d}(f)$ for arbitrarily large $n$ because they could not enjoy two different routes forwards and backwards through the resulting bottleneck in the hyperplane $x_{1}=x_{0}$. In discussing (2.8) we therefore impose, without essential loss of generality, the condition

$$
\begin{equation*}
f_{d}(x) \geqslant 1 \tag{2.9}
\end{equation*}
$$

We shall prove that (2.6) and (2.9) are sufficient conditions for the truth of

$$
\begin{equation*}
\liminf _{n \rightarrow \infty} n^{-1} \log p_{n}(f) \geqslant \kappa \tag{2.10}
\end{equation*}
$$

and hence for the truth of (2.5) and (2.8) in view of the trivial inequality

$$
\begin{equation*}
p_{n}(f) \leqslant c_{n-1}(f) \leqslant c_{n-1} \tag{2.11}
\end{equation*}
$$

coupled with (2.4). (Note that (2.9) is no longer needed for (2.5), because a separate return path to the origin is then unnecessary.)

## 3. Bellman's theorem

A pattern $P$ is a prescribed finite sequence of steps, and is said to occur on a walk $w$ if it is a subwalk of $w$. Successive occurrences may overlap: for example, when $d=2$, we may write $w=\mathrm{NE}^{3} \mathrm{SE}^{2} \mathrm{~N}$ for the walk whose successive steps are north, east, east, east, south, east, east, north and the pattern $\mathrm{P}=\mathrm{E}^{2}$ occurs three times on $w$. Given any particular pattern $P$, we can always find a Pólya walk on which $P$ occurs as often as we wish; but the situation is quite different if $w$ is a SAw. Obviously P cannot occur on any sAw unless $P$ is a saw itself. It is also possible to find patterns $P$ that may occur once or twice on saws in $\mathbb{Z}^{d}$ though never more often than twice. Thus the 12-stepped pattern $\mathrm{P}=\mathrm{NWS}^{2} E^{4} \mathrm{~N}^{2} \mathrm{WS}$ occurs twice on the 35 -sAw $w=P W N^{2} E^{5} S^{2} \mathrm{WP}$, but $P$ can never occur thrice on any SAW in $\mathbb{Z}^{2}$; and this sort of 'lobster-pot' or 'trapdoor' construction can readily be generalised to $\mathbb{Z}^{d}$ for any $d \geqslant 2$. On the other hand, if $k>2$, there are no patterns that can occur $k$ times on a saw but never $k+1$ times, for the following reason.

Theorem. If a particular pattern P occurs three times on some saw in $\mathbb{Z}^{d}$, then P can occur infinitely often on an infinite saw in $\mathbb{Z}^{d}$.

This theorem, intuitively motivated by the fact that a walk has just two ends, needs careful proof because successive occurrences of P may overlap and become highly entangled with each other in the rich topology of high-dimensional space.

Let $P_{1}, P_{2}, P_{3}$ be the first, second and third occurrences of $P$ on a sAw $w$; and write $B_{1}, B_{2}, B_{3}$ for the respective snug boxes of $P_{1}, P_{2}, P_{3}$. These three boxes all have the same shape, size, and orientation, because the same pattern $P$ is snug in each; but $B_{1}$, $B_{2}, B_{3}$ must occupy different positions in $\mathbb{R}^{d}$, for $B_{1}=B_{2}$ would imply that $P_{1}$ and $P_{2}$ exactly coincided in violation of the self-avoidance of $w$. Hence at least one complete face of $B_{1}$ must lie strictly outside $B_{2}$. But $P_{1}$, being snug in $B_{1}$, visits every face of $B_{1}$; so $P_{1}$ contains a vertex outside $B_{2}$. This vertex cannot lie in any overlap of $P_{1}$ and $P_{2}$, because $P_{2}$ lies wholly in $B_{2}$; so this vertex precedes all vertices of $P_{2}$ on $w$. Hence there is a non-empty class of vertices on $w$ lying outside $\mathrm{B}_{2}$ and preceding all vertices of $P_{2}$; and this class possesses a latest vertex $V_{1}$ on $w$. Similar consideration of $B_{2} \neq B_{3}$ guarantees the existence of an earliest vertex $V_{3}$ outside $B_{2}$ and succeeding all vertices of $P_{2}$. Let $w_{2}$ denote the subwalk of $w$ from $V_{1}$ to $V_{3}$. Discard all steps of $w$ except the steps of $w_{2}$. Then $w_{2}$ has just two vertices outside $B_{2}$, and there exists a box $B_{4}$ (whose interior strictly contains $\mathrm{B}_{2}$ ) such that only the first and last vertices of $W_{2}$ lie on the faces of $\mathrm{B}_{4}$. ( $\mathrm{B}_{4}$ is not snug for $w_{2}$.) If $\mathrm{V}_{1}$ and $\mathrm{V}_{3}$ do not lie on opposite faces of $B_{4}$, we may append some extra steps to the end of $w_{2}$ to make this happen; and we can ensure that the extended walk is still a saw by suitably utilising steps confined to the surface of $B_{4}$. Finally we can add one more step to the end of this extended $w_{2}$ to reach a vertex outside $B_{4}$. This yields a pattern $Q$, which is a sAw containing the pattern $P$, such that the infinite repetition QQQ ... is an infinite sAw on which $P$ occurs infinitely often.

This theorem is known as Bellman's theorem on account of the Bellman's remark 'What I tell you three times is true' (Carroll 1876); but Carroll needed 42 boxes 'all carefully packed' in his account.

## 4. Proof of (2.10)

Prescribe an even integer $N>0$ and consider the set of points of $\mathbb{Z}^{d}$ satisfying

$$
\begin{equation*}
x_{1}=0, \quad 0 \leqslant x_{k} \leqslant N \quad(2 \leqslant k \leqslant d) \tag{4.1}
\end{equation*}
$$

This set contains $M=(d-1) N^{d-1}$ unordered pairs of adjacent points, which we index in some fixed manner by $i=1,2, \ldots, M$. Now consider an $N$-sap translated so that its snug box has the form

$$
\begin{equation*}
0 \leqslant x \leqslant b \tag{4.2}
\end{equation*}
$$

No coordinate of $b$ can exceed $N$, and the $N$-sap must pass through at least one pair of adjacent vertices on the face $x_{1}=0$. So we can index this $N$-sap by the index $i$ ( $1 \leqslant i \leqslant M$ ), using the smallest available index if the $N$-sap passes through more than one pair of adjacent vertices on the face $x_{1}=0$. Similarly, we can classify the $N$-sap by a second index $j(1 \leqslant j \leqslant M)$, corresponding to a pair of adjacent vertices that it visits on the face $x_{1}=b_{1}$ of its snug box (using, of course, the same fixed indexing as in (4.1) but with $x_{1}=b_{1}$ ). The indexing $i, j$ partitions the set of all inequivalent $N$-saps into $M^{2}$ disjoint subsets $P_{i}$, of which some subsets may be empty.

Now consider two $N$-saps, the first a member of $P_{i j}$ fitted into a snug box (4.2), and the second a member of $P_{j k}$ translated to fit into a snug box

$$
\begin{equation*}
b_{1}+1 \leqslant x_{1} \leqslant b_{1}^{\prime}, \quad 0 \leqslant x_{k} \leqslant b_{k}^{\prime} \quad(2 \leqslant k \leqslant d) . \tag{4.3}
\end{equation*}
$$

Since the second suffix of $P_{i j}$ equals the first suffix of $P_{j k}$, the two saps will have parallel adjacent edges on the faces $x_{1}=b_{1}$ of the first SAP and $x_{1}=b_{1}+1$ of the second SAP. We can remove these two edges and replace them by a pair of edges between the two faces, thus concatenating the two saps and constructing a $2 N$-sap. Further, given any 2 N -sap constructed in this manner, we can reverse the process and uniquely recover the pair of $N$-saps from which it was derived. Hence, if $p_{i j}$ denotes the number of $N$-saps in the subset $P_{i j}$, and $\boldsymbol{P}$ is the $M \times M$ matrix $\boldsymbol{P}=\left(p_{i j}\right)$, the total number of $2 N$-saps constructed by this concatenation and belonging to the subset $P_{j k}$ is the $(j, k)$-element in $\boldsymbol{P}^{2}$. By repeated concatenation, we can construct a set of inequivalent Nr -saps, whose number is the sum of the elements in the matrix $\boldsymbol{P}^{r}$. Moreover all these $N r$-saps can be translated to fit into snug boxes of the form

$$
\begin{equation*}
a_{1} \leqslant x_{1} \leqslant b_{1}, \quad 0 \leqslant x_{k} \leqslant b_{k} \leqslant N \quad(2 \leqslant k \leqslant d) \tag{4.4}
\end{equation*}
$$

and all these snug boxes will be contained in $\mathbb{Z}^{d}(f)$ provided that

$$
\begin{equation*}
a_{1} \geqslant \alpha(N) \tag{4.5}
\end{equation*}
$$

for some number $\alpha$ depending on $N$, in view of (2.6).
We now wish to construct a polygon, in $\mathbb{Z}^{d}(f)$, which visits the origin and contains the edge between the points $\left(u_{1}, u_{2}, \ldots, u_{q}, \ldots, u_{d}\right)$ and ( $u_{1}, u_{2}, \ldots, u_{q}+1, \ldots, u_{d}$ ) in the face $x_{1}=u_{1}=b_{1}$ of its snug box $\mathbf{0} \leqslant \boldsymbol{x} \leqslant \boldsymbol{b}$. If we write $\boldsymbol{e}_{k}^{u}$ for a succession of $u$ steps in the direction $\boldsymbol{e}_{k}$ and $\tilde{e}_{k}^{u}$ for a succession of $u$ steps in the direction $-\boldsymbol{e}_{k}$, a polygon satisfying these conditions can be written as

$$
\Pi_{1}=\boldsymbol{e}_{1}^{u_{1}-1} \boldsymbol{e}_{2}^{u_{2}} \ldots \boldsymbol{e}_{d}^{u_{d}} \boldsymbol{e}_{1} \boldsymbol{e}_{q} \bar{e}_{d} \overline{\boldsymbol{e}}_{1} \bar{e}_{q} \bar{e}_{d}^{u_{d}-2} \overline{\boldsymbol{e}}_{d-1}^{u_{d-1}} \ldots \overline{\boldsymbol{e}}_{2}^{u_{2}} \bar{e}_{1}^{u_{1}-1}
$$

provided that $u_{d} \geqslant 2$. If $u_{d}=0$ or 1 , an appropriate polygon is

$$
\Pi_{2}=\boldsymbol{e}_{1}^{u_{1}-1} \boldsymbol{e}_{2}^{u_{2}} \ldots e_{d}^{u_{d}} e_{1} e_{q} e_{d} \bar{e}_{1} \bar{e}_{q} \bar{e}_{d} u_{d} \overline{\boldsymbol{e}}_{d-1}^{u_{d-1}} \ldots \overline{\boldsymbol{e}}_{2}^{u_{2}} \overline{\boldsymbol{e}}_{1}^{u_{1}-1} .
$$

After translation, $\Pi_{1}$ fits into the snug box

$$
\begin{equation*}
0 \leqslant x_{k} \leqslant u_{k}+\delta_{k} \tag{4.6}
\end{equation*}
$$

with

$$
\delta_{k}= \begin{cases}1 & \text { if } k=q  \tag{4.7}\\ 0 & \text { otherwise }\end{cases}
$$

and $\Pi_{2}$ fits into the snug box (4.6) with

$$
\delta_{k}= \begin{cases}1 & \text { if } k=q \text { or } d  \tag{4.8}\\ 0 & \text { otherwise }\end{cases}
$$

These sAPs will belong to some class $P_{i j}$ provided that $u_{k}<N$ for $2 \leqslant k \leqslant d$. By choosing an integer $q(2 \leqslant q \leqslant d)$ and integers $u_{k}<N, k=2, \ldots, d$, we can construct for some $i_{0}$ a SAP belonging to $P_{i_{0}}$ for any given value of $j=1,2, \ldots, M$, and this polygon can be concatenated in front of any $N r$-SAP by taking $a_{1}=u_{1}+1$. The SAP formed by this concatenation will be contained in $\mathbb{Z}^{d}(f)$ provided that

$$
\begin{equation*}
u_{1} \geqslant \beta(N) \tag{4.9}
\end{equation*}
$$

for some integer $\beta(N)$, because of (2.6). By choosing $u_{1}$ appropriately, subject to (4.9) the SAP can be constructed with $n$ edges where $n$ is any sufficiently large even integer. This $n$ will have the form

$$
\begin{equation*}
n=N r+A \tag{4.10}
\end{equation*}
$$

where $A$ is the number of edges in $\Pi_{1}$ or $\Pi_{2}$, and $r$ is any positive integer. Hence

$$
\begin{equation*}
p_{n}(f) \geqslant \text { sum of elements of } \boldsymbol{P}^{r} \geqslant p_{j k} p_{k j} p_{j k} \ldots \quad(r \text { terms }) \tag{4.11}
\end{equation*}
$$

for any pair of indices $j, k$. By reflecting any $N$-saP in $x_{1}=0$ it is easy to see that $P$ is a symmetric matrix. If we choose $j, k$ so that $p_{j k}$ is the largest element in $\boldsymbol{P}$ we have

$$
\begin{equation*}
p_{n}(f) \geqslant\left(p_{N} / M^{2}\right)^{r} \tag{4.12}
\end{equation*}
$$

because $p_{N}$ is the sum of all elements in $\boldsymbol{P}$. Hence
$n^{-1} \log p_{n}(f) \geqslant N^{-1}(1-A / n) \log p_{N}-2 N^{-1}(1-A / n) \log \left[(d-1) N^{d-1}\right]$.
Letting $n \rightarrow \infty$ in (4.13) we have

$$
\begin{equation*}
\liminf _{n \rightarrow \infty} n^{-1} \log p_{n}(f) \geqslant N^{-1}\left\{\log p_{N}-2 \log \left[(d-1) N^{d-1}\right]\right\} . \tag{4.14}
\end{equation*}
$$

Since the lhs of (4.14) is independent of $N$, we can now let $N \rightarrow \infty$ and (2.10) follows from (4.14) and (2.8). The influence of $f$ upon the rate of approach to the limit can be estimated from (4.13).

## 5. Constrained self-avoiding walks and patterns

In $\S 4$ we discussed a sufficient condition on the functions $f_{k}$ defining an $f$-wedge which ensured that the connective constant of saws in an $f$-wedge would be identical to that for saws on the complete lattice. We now turn to the opposite case and look for sufficient conditions for there to be exponentially fewer saws in the $f$-wedge than on the complete lattice. The idea is to make use of a theorem of Kesten (1963) on patterns.

If $\mathbf{P}$ is any (finite) pattern and $c_{n}(\varepsilon, \mathbf{P})$ is the number of $n$-step self-avoiding walks in which the pattern P occurs at most $\varepsilon n$ times then there exists a value of $\varepsilon>0$ such that

$$
\begin{equation*}
\limsup _{n \rightarrow \infty}\left[c_{n}(\varepsilon, P) / c_{n}\right]^{1 / n}<1 \tag{5.1}
\end{equation*}
$$

provided that there exists a self-avoiding walk in which P appears more than twice.
The corollary which we require is as follows: if there exists a pattern P which occurs more than twice in at least one self-avoiding walk in $\mathbb{Z}^{d}$ but does not occur in any self-avoiding walk in $\mathbb{Z}^{d}(f)$ then

$$
\begin{equation*}
\limsup _{n \rightarrow \infty} n^{-1} \log c_{n}(f)<\kappa \tag{5.2}
\end{equation*}
$$

Since $c_{n}(f) \leqslant c_{n}(0, \mathrm{P})$ (5.2) follows immediately from (5.1).
We now consider two examples.
Example 1. Consider a saw on $\mathbb{Z}^{d}$ confined to lie in or between the hyperplanes $x_{d}=0$ and $x_{d}=L$. Let $c_{n}(L)$ be the number of these walks which start at the origin. If N and S represent steps in the positive $x_{d}$ and negative $x_{d}$ directions, and E and W represent steps in the positive $x_{1}$ and negative $x_{1}$ directions, then the pattern

$$
\mathrm{P}=\mathrm{EN}^{L+1} \mathrm{ES}^{L+1}
$$

does not occur in any saw confined between $x_{d}=0$ and $x_{d}=L$. However, P does occur more than twice in a SAW in $\mathbb{Z}^{d}$ (the walk $\mathbf{P}^{3}$ for example) so that

$$
\begin{equation*}
\limsup _{n \rightarrow \infty} n^{-1} \log c_{n}(L)<\kappa \tag{5.3}
\end{equation*}
$$

for any finite $L$.
Example 2. Consider a sAw on $\mathbb{Z}^{2}(f)$ with $f_{2}(x)=\mathrm{e}^{x} \sin ^{2}(\pi x / L) . f_{2}(x)$ is unbounded above but is zero when $x$ is an integral multiple of $L$. The pattern

$$
\mathrm{P}=\mathrm{E}^{L+1} \mathrm{NE}^{L+1}
$$

occurs more than twice in a SAw in $\mathbb{Z}^{2}$ (e.g. in the walk $P^{3}$ ) but does not occur in any saw in $\mathbb{Z}^{2}(f)$. Hence

$$
\begin{equation*}
\limsup _{n \rightarrow \infty} n^{-1} \log c_{n}(f)<\kappa \tag{5.4}
\end{equation*}
$$

A similar argument can easily be constructed for any function $f_{2}$ which periodically takes any prescribed finite value. However, if $f_{2}$ only returns to the prescribed finite value aperiodically and sufficiently increasingly rarely, then (5.4) could be false.

## 6. Walks between two parallel planes

In this section we consider the subset $C(n, L) \subseteq C(n)$ such that the walks lie in or between the two hyperplanes $x_{d}=0$ and $x_{d}=L$. If we write $x_{k}(l)$ for the $k$ th coordinate of the $l$ th vertex in the walk, a walk is a member of $C(n, L)$ if

$$
\begin{equation*}
0 \leqslant x_{d}(l) \leqslant L, \quad l=0,1, \ldots, n . \tag{6.1}
\end{equation*}
$$

It will be convenient to construct a subset $B(n, L) \subseteq C(n, L)$ such that

$$
x_{1}(0)<x_{1}(l) \leqslant x_{1}(n) \quad l=1,2, \ldots, n
$$

and

$$
\begin{equation*}
x_{d}(n)=0 . \tag{6.2}
\end{equation*}
$$

If the numbers of members of $B(n, L)$ and $C(N, L)$ are $b_{n}(L)$ and $c_{n}(L)$ respectively it is straightforward (Whittington 1983) to show that
$0<\sup _{n>0} n^{-1} \log b_{n}(L)=\lim _{n \rightarrow \infty} n^{-1} \log b_{n}(L)=\lim _{n \rightarrow \infty} n^{-1} \log c_{n}(L) \equiv \kappa(L)<\infty$.
The arguments of $\S 5$ establish that

$$
\begin{equation*}
\kappa(L)<\kappa \tag{6.4}
\end{equation*}
$$

for any finite $L$. Here we shall prove that $\kappa(L)$ is a strictly monotone increasing function of $L$ and that

$$
\begin{equation*}
\lim _{L \rightarrow \infty} \kappa(L)=\kappa . \tag{6.5}
\end{equation*}
$$

We define the subset $P(n, L) \subseteq B(n, L)$ such that a walk in $P(n, L)$ cannot be decomposed into two walks, one from $B(m, L)$ and one from $B(n-m, L)$, for any $m$. For this reason we call $P(n, L)$ the set of prime walks. If $p_{n}(L)$ is the number of members of $P(n, L)$ we can write

$$
\begin{equation*}
b_{n}(L)=\sum_{m=1}^{n} p_{m}(L) b_{n-m}(L) \tag{6.6}
\end{equation*}
$$

and if we form generating functions

$$
\begin{equation*}
\mathscr{P}(x, L)=\sum_{m} p_{m}(L) x^{m} \tag{6.7}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathscr{B}(x, L)=\sum_{m} b_{m}(L) x^{m}, \tag{6.8}
\end{equation*}
$$

we obtain

$$
\begin{equation*}
\mathscr{B}(x, L)=\mathscr{P}(x, L) /[1-\mathscr{P}(x, L)] . \tag{6.9}
\end{equation*}
$$

We now remark that $\mathscr{B}(x, L) \rightarrow \infty$ as $x \rightarrow \exp [-\kappa(L)]$ from below. Indeed this follows by an argument exactly analogous to an argument of Kesten (1963, theorem 5), and we shall not repeat the details here. This, together with (6.9), implies that

$$
\begin{equation*}
\mathscr{P}\left(\mathrm{e}^{-\kappa(L)}, L\right)=1 \tag{6.10}
\end{equation*}
$$

But if $\kappa(L)=\kappa\left(L^{\prime}\right), L>L^{\prime}$, then $\mathscr{P}(x, L)=1$ and $\mathscr{P}\left(x, L^{\prime}\right)=1$ must have a common root at $x=\mathrm{e}^{-\kappa(L)}$. This is impossible since $p_{m}(L)$ is non-negative and $p_{m}(L) \geqslant p_{m}\left(L^{\prime}\right)$ with strict inequality for some values of $m$. Hence

$$
\begin{equation*}
\kappa(L)>\kappa\left(L^{\prime}\right), \quad L>L^{\prime} . \tag{6.11}
\end{equation*}
$$

If we put $f_{d}(x)=L$ for $x \geqslant 0$, (4.14) becomes

$$
\begin{align*}
\kappa & \geqslant \liminf _{n \rightarrow \infty} n^{-1} \log c_{n}(L) \\
& \geqslant \liminf _{n \rightarrow \infty} n^{-1} \log p_{n}(L) \\
& \geqslant L^{-1}\left\{\log p_{L}-2 \log \left[(d-1) L^{d-1}\right]\right\} \tag{6.12}
\end{align*}
$$

The Rhs of (6.12) tends to $\kappa$ as $L \rightarrow \infty$. Hence

$$
\begin{equation*}
\lim _{L \rightarrow \infty} \kappa(L)=\kappa . \tag{6.13}
\end{equation*}
$$

## 7. Discussion

We have been concerned with the number of self-avoiding walks confined to lie in a region between the $x$ axis and a curve $y=f(x)$, and the generalisation of this restriction to $d$ dimensions. Provided that $f \rightarrow \infty$ as $x \rightarrow \infty$ we have shown that the connective constant for such walks is equal to the connective constant for the complete lattice. In particular, this establishes that the three kinds of walks in a wedge, corresponding to the zero spin space dimension limit of Cardy's problem (Cardy 1983) have the same connective constant as the complete lattice, for all positive wedge angles.

In § 5 we discussed the use of a theorem by Kesten, on patterns in self-avoiding walks, to establish conditions on $f$ sufficient to ensure that only an exponentially small fraction of self-avoiding walks survive the condition that they are restricted to an $f$-wedge, and in $\S 6$ we discussed the particular case of walks confined between parallel lines or planes, in more detail.

The result derived in $\S 4$ should be contrasted with a result of Grimmett (1983) for the bond percolation problem. Grimmett investigated the percolation probability on the subset of the square lattice $\left\{(x, y) \in \mathbb{Z}^{2}: 0 \leqslant y \leqslant f(x), 0 \leqslant x\right\}$. If $f(x)=a \log (x+1)$ he showed that the critical probability for percolation $p_{\mathrm{c}}(a)$ is a continuous strictly decreasing function of $a$ with $p_{\mathrm{c}}(a) \rightarrow 1$ as $a \rightarrow 0+$ and $p_{\mathrm{c}}(a) \rightarrow \frac{1}{2}$, as $a \rightarrow \infty$. Hence, if $f(x) / \log x \rightarrow \infty$ as $x \rightarrow \infty, p_{c}=\frac{1}{2}$ and if $f(x)=o(\log x), p_{c}=1$. By contrast, in the selfavoiding walk problem the connective constant is independent of $f$ if $f \rightarrow \infty$ no matter how slowly.

Grimmett's arguments depend strongly on the lattice being two-dimensional, but weaker results can be obtained for $d>2$. We consider $\mathbb{Z}^{d}(f)$ and let $g(x)$ be the number of bonds in the $x_{1}$ direction from $x_{1}=x$ to $x_{1}=x+1$ in $\mathbb{Z}^{d}(f)$. Let $A_{x}$ be the event that all of these bonds are blocked, and $q$ be the probability that a given bond is blocked. Since the events $\left\{A_{x}\right\}$ are independent it follows from the Borel-Cantelli lemma that infinitely many $A_{x}$ will occur with probability one if

$$
\begin{equation*}
\sum_{x} q^{g(x)}=\infty \tag{7.1}
\end{equation*}
$$

If, for $x \geqslant x_{0}, g(x) \leqslant a \log x$ then

$$
\begin{equation*}
\sum_{x \geqslant x_{0}} q^{g(x)} \geqslant \sum_{x \geqslant x_{0}} x^{a \log q} \tag{7.2}
\end{equation*}
$$

which is infinite provided that $a \log q \geqslant-1$. In this case no percolation occurs and
hence we have shown that

$$
\begin{equation*}
p_{c}(f) \geqslant 1-\mathrm{e}^{-1 / a} . \tag{7.3}
\end{equation*}
$$

In particular, if $g(x)=o(\log x)$ then $p_{c}=1$.
We can also show that $p_{c}<1$ if $f_{k}(x) \geqslant \varepsilon(\log x)^{1 / d}$ for $k \geqslant 2$, for $x \geqslant x_{0}$, for any positive $\varepsilon$. We sketch the argument only for $d=3$. If we consider a plane $x_{1}=x$, we , can order the $g(x)$ vertices in this plane which are also in $\mathbb{Z}^{3}(f)$ as $(0,0),(0,1),(1,1)$, $(1,0),(2,0),(2,1),(2,2),(1,2),(0,2),(0,3),(1,3), \ldots$ where we list only the $\left(x_{2}, x_{3}\right)$ coordinates. We now consider the sublattice consisting of these vertices for all $x_{1}$ planes, the bonds between adjacent $x_{1}$ planes and bonds within planes which join adjacent members in the foregoing ordering. By the containment theorem (Fisher 1961) this lattice percolates less readily than $\mathbb{Z}^{3}(f)$. However, this lattice can be deformed into a subset of the two-dimensional square lattice ('unroll' the vertices in the ordered set above) so we can make use of Grimmett's result (that $p_{c}(a)$ is a strictly decreasing function of $a$ ) to establish that $p_{c}<1$ provided that $f_{k}(x) \geqslant \varepsilon(\log x)^{1 / 2}$, $\varepsilon>0, x \geqslant x_{0}, k=2,3$. Likewise

$$
\begin{equation*}
f_{k}(x) /(\log x)^{1 / 2} \rightarrow \infty \quad \text { as } x \rightarrow \infty \quad(k=2,3) \tag{7.4}
\end{equation*}
$$

implies $p_{\mathrm{c}}(f) \leqslant \frac{1}{2}$. Under additional conditions, it might be true that (7.4) implies that $p_{c}(f)=p_{\mathrm{c}}\left(\mathrm{i}\right.$. . the critical probability of $\left.\mathbb{Z}^{3}\right)$. Grimmett (1984) has verified our conjecture that (7.4) could be replaced by

$$
\begin{equation*}
g(x) / \log x \rightarrow \infty \quad \text { as } x \rightarrow \infty \tag{7.5}
\end{equation*}
$$

We can now ask for corresponding results for the Ising problem. If we consider the Ising problem on an infinite subset of the square lattice, how does the critical temperature depend on the subset chosen? There are a few results on the 'slab' geometry discussed in §6. If $d=2$ there is no phase transition, on any finite width slab, but if $d=3$ there is series analysis work on the dependence of the critical temperature ( $T_{\mathrm{c}}$ ) on the width ( $L$ ) (Allan 1970), though we are not aware of any rigorous results on this question. However, the spherical model has been solved in a slab geometry and the dependence of $T_{c}$ on $L$ is known for this problem (see e.g. Watson 1972).

For the Ising problem on an $f$-wedge, even in two dimensions, almost nothing seems to be known. It is tempting to conjecture that the behaviour follows the percolation pattern rather than the self-avoiding walk pattern and it would be interesting to investigate this. The only result of which we are aware is that the half-plane and the plane square lattices have the same critical temperature (McCoy and Wu 1973). On the other hand, the critical exponents of the correlation function and of the spontaneous magnetisation do depend upon the wedge angle (Barber et al 1984).

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